## QFT II - PROBLEM SET 2

## (29) CROSS SECTIONS FOR SIMPLE CASES

The cross section for two particles in the final state is given by

$$\left(\frac{d\sigma}{d\Omega}\right)_{cm} = \frac{1}{2E_A 2E_B |v_A - v_B|} \frac{|\mathbf{p_1}|}{(2\pi)^2 4E_{cm}} |\mathcal{M}(p_A, p_B \to p_1, p_2)|^2,$$

where the incoming particles are labeled A, B, the outgoing particles  $p_1, p_2$  and  $E_{cm}$  is the center of mass energy.

a) Suppose that one of them (particle A) is massless, like in the left figure below. Denote  $p_A^{\mu} = (k, k\hat{z})$ ,  $p_B^{\mu} = (E, -k\hat{z})$ . What is the cross section in this case? *Hint: use*  $\boldsymbol{p} = m\boldsymbol{v}/\sqrt{1-v^2}$  to get an expression for  $v_B$  in terms of  $\boldsymbol{p}_B = -k\hat{z}$  and  $E_B = E$ .

b) Suppose that all particles have the same mass, like in the right figure below. What is the cross section then?



## (30) FEYNMAN RULES FOR EVERYONE

In general, you can derive Feynman rules for any theory by simply taking derivatives of the action S with respect to the fields. The momentum conserving  $\delta$  functions that occur in the derivatives are omitted in the expressions for the propagator and the vertex. They do, however yield the rules for momentum conservation at each vertex.

Drawing a Propagator in a Feynam diagram is associated with

$$P = \frac{i}{\mathcal{S}^{(2)}},$$

where

$$\mathcal{S}^{(2)} = \frac{\delta^2 S(\phi, \chi, \dots)}{\delta \phi \delta \chi}_{|\phi = \chi = \dots = 0}.$$

Likewise, the vertices are higher derivatives of S, again evaluated at vanishing field values. For instance the vertex for the interaction Lagrangian  $\mathcal{L}(\phi, \chi) = \lambda \chi \chi \phi \phi / 4$  is

$$V = i \frac{\delta^4 S(\phi, \chi, \dots)}{\delta^2 \phi \delta^2 \chi}_{|\phi = \chi = \dots = 0} = i \lambda$$

a) To flex our muscles a bit, let's apply this for several theories. So what are the propagator(s) and vertex (vertices) for the following Lagrangians ?

i) Scalar  $\phi^4$  theory

$$\mathcal{S} = -\int_{p} \frac{1}{2} (p_{\mu}p^{\mu} + m^{2})\phi(p)\phi(-p) - \frac{\lambda}{4!} \int_{p_{1}...p_{4}} \delta(p_{1} + p_{2} + p_{3} + p_{4})\phi(p_{1})\phi(p_{2})\phi(p_{3})\phi(p_{4})$$

ii) Complex scalar  $\phi^4$  theory

$$S = -\int_{p} (p_{\mu}p^{\mu} + m^{2})\phi^{*}(p)\phi(p) - \frac{\lambda}{4}\int_{p_{1}\dots p_{4}} \delta(p_{1} - p_{2} + p_{3} - p_{4})\phi(p_{1})\phi^{*}(p_{2})\phi(p_{3})\phi^{*}(p_{4})$$

b) The interaction of a dipole with the electric field is described by the following action

$$S_{int} = -\int_x \frac{i}{2}\bar{\Psi}(x)\sigma^{\mu\nu}(M+\gamma^5 D)\Psi(x)F_{\mu\nu}(x),$$

where  $\sigma^{\mu\nu} = \frac{i}{2} [\gamma^{\mu} \gamma^{\nu} - \gamma^{\nu} \gamma^{\mu}]$  is the usual spin operator and M and D are the magnetic and electric dipole strength.

(i) Fourier transform  $S_{int}$  to momentum space (please use the letter k for the momentum of  $A_{\mu}$ ).

(ii)  $\sigma_{\mu\nu}$  is antisymmetric in  $\mu, \nu$ . Use this to add a "clever" 0 to  $F_{\mu\nu}$  to show that you can simplify  $\sigma^{\mu\nu}F_{\mu\nu} \rightarrow 2\sigma^{\mu\nu}k_{\mu}A_{\nu}$ , where k is the photon momentum.

(iii) What is the vertex for this interaction ? Draw it. What happens to the vertex for very soft photons ?

(iv) Draw the Feynman diagramm for the scattering of an electron and a dark matter fermion that carries an electric dipole moment D but no charge.

(v) Obtain the amplitude  $\mathcal{M}$  for this process.

## (31) BHABHA SCATTERING

Let us consider electron-positron (Bhaba) scattering. At tree level, there are two diagrams contributing. One in which electrons and positrons simply scatter and one where they annihilate and re-appear later.

a) Draw the Feynman diagrams for both processes.

b) Use the QED Feynman rules to obtain the amplitude  $\mathcal{M}$  and  $|\mathcal{M}|^2$ .

c) Average over incoming spins and sum over outgoing spins, i.e. use the spin sum rules to obtain  $|\mathcal{M}|^2$  in terms of Traces.

d) Stop here. Don't evaluate the traces ...