Examining nonextensive statistics in relativistic heavy-ion collisions

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Simon, A., & Wolschin, G. (2018). Examining nonextensive statistics in relativistic heavy-ion collisions. *Physical Review C*, 97(4), 044913.

Setting and Problem

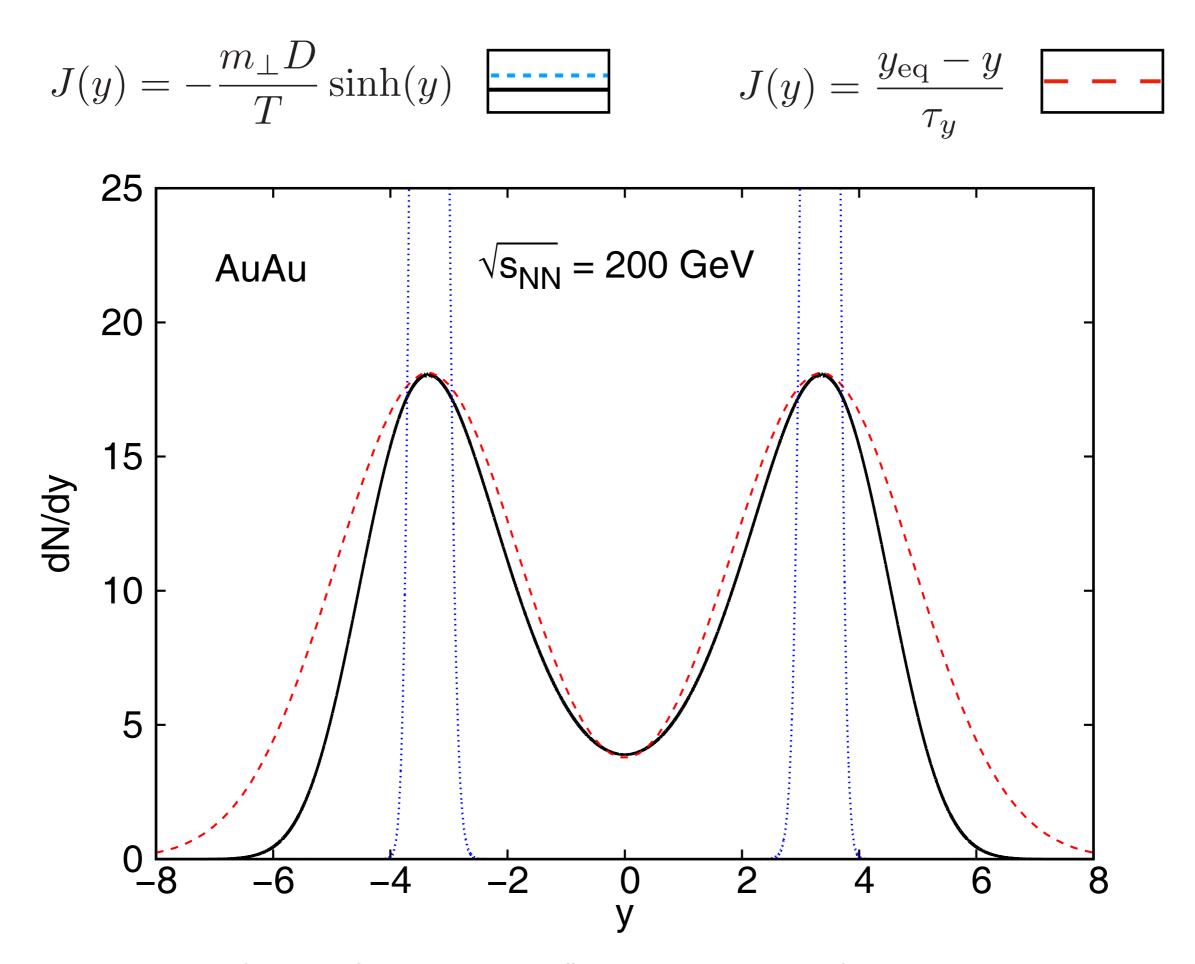
- Model stopping of (proton antiproton) @ heavy-ion collisions (RHIC, PbPb)
- Use relativistic diffusion model

$$\frac{\partial}{\partial t}W(y,t) = -\frac{\partial}{\partial y}[J(y,t)W(y,t)] + \frac{\partial^2}{\partial y^2}[D(y,t)W(y,t)]$$

$$- \text{drift} - \text{diffusion}$$

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- Know drift coefficient J(y,t) by:
 - interaction time and peak position: <u>amplitude</u>
 - stationary solution: <u>functional form</u>
- Determine diffusion by fluctuation-dissipation theorem



Forndran, F., & Wolschin, G. (2017). Relativistic diffusion model with nonlinear drift. The European Physical Journal A, 53(2), 37.

Account for collective expansion

- 'Artificially' through stronger diffusion coefficient, or
- Change underlying equation: <u>non-linear Fokker-Planck</u> <u>equation</u>

$$\frac{\partial}{\partial t}W(y,t)^{\mu} = -\frac{\partial}{\partial y}[J(y,t)W(y,t)^{\mu}] + \frac{\partial^2}{\partial y^2}[D(y,t)W(y,t)^{\nu}]$$

Non-extensive entropy

Boltzmann-Gibbs entropy

q entropy

$$S_{BG} = -k\sum_{i} p_i \ln p_i$$

 $S_B(\mathbf{A} + \mathbf{B}) = S_B(\mathbf{A}) + S_B(\mathbf{B})$

$$S_q = \frac{k}{q-1} \left(1 - \sum_i p_i^q\right)$$

 $S_q(\mathbf{A} + \mathbf{B}) = S_q(\mathbf{A}) + S_q(\mathbf{B})$ $+ (1 - q)S_q(\mathbf{A})S_q(\mathbf{B})$

 $S_q = \left\langle \ln_q \frac{1}{n_i} \right\rangle$

Connection to NLFPE?

• Maximize entropy
$$S_q[p] = rac{1 - \int du[p(u)]^q}{q - 1}$$

$$p_q(y,t) = \frac{\{1 - \beta(t)(1 - q)[y - y_m(t)]^2\}^{1/(1-q)}}{Z_q(t)}$$

- Insert as ansatz into NLFPE and chose linear drift
- Leads to: $q = 1 + \mu v$

Tsallis, C., & Bukman, D. J. (1996). Anomalous diffusion in the presence of external forces: Exact time-dependent solutions and their thermostatistical basis. *Physical Review E*, *54*(3), R2197.

Final equation

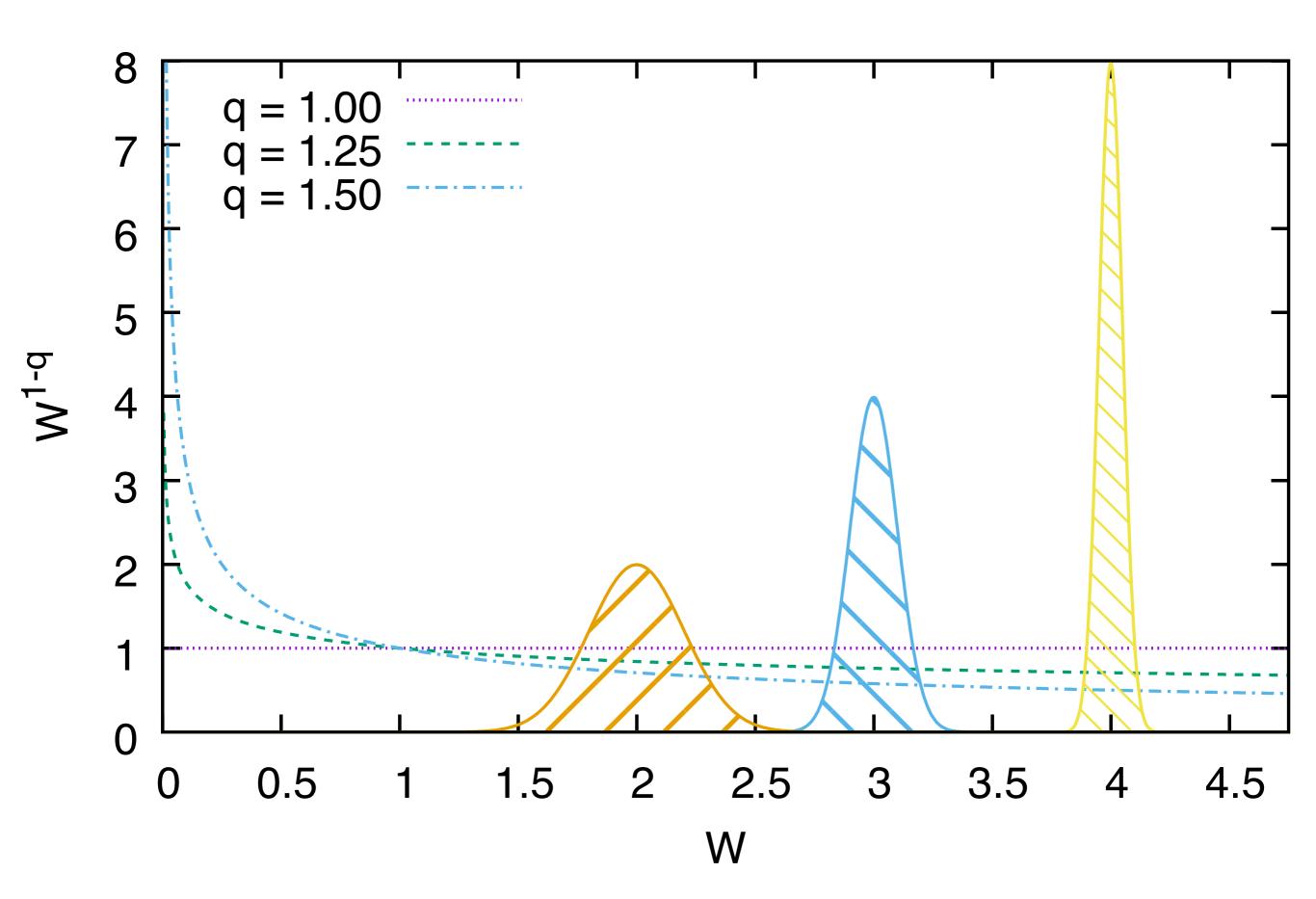
• Rescale time and get

$$\frac{\partial f}{\partial \tau} = \frac{\partial}{\partial y} [\sinh(y) f(y,t)] + \gamma \frac{\partial^2}{\partial y^2} [f(y,t)^{2-q}], \ \gamma = T/m_{\perp}$$

• Non-linear diffusion coefficient

$$\frac{\partial^2}{\partial y^2} [DW^{2-q}] = \frac{\partial^2}{\partial y^2} [(DW^{1-q})W]$$

• What changes?



FEM

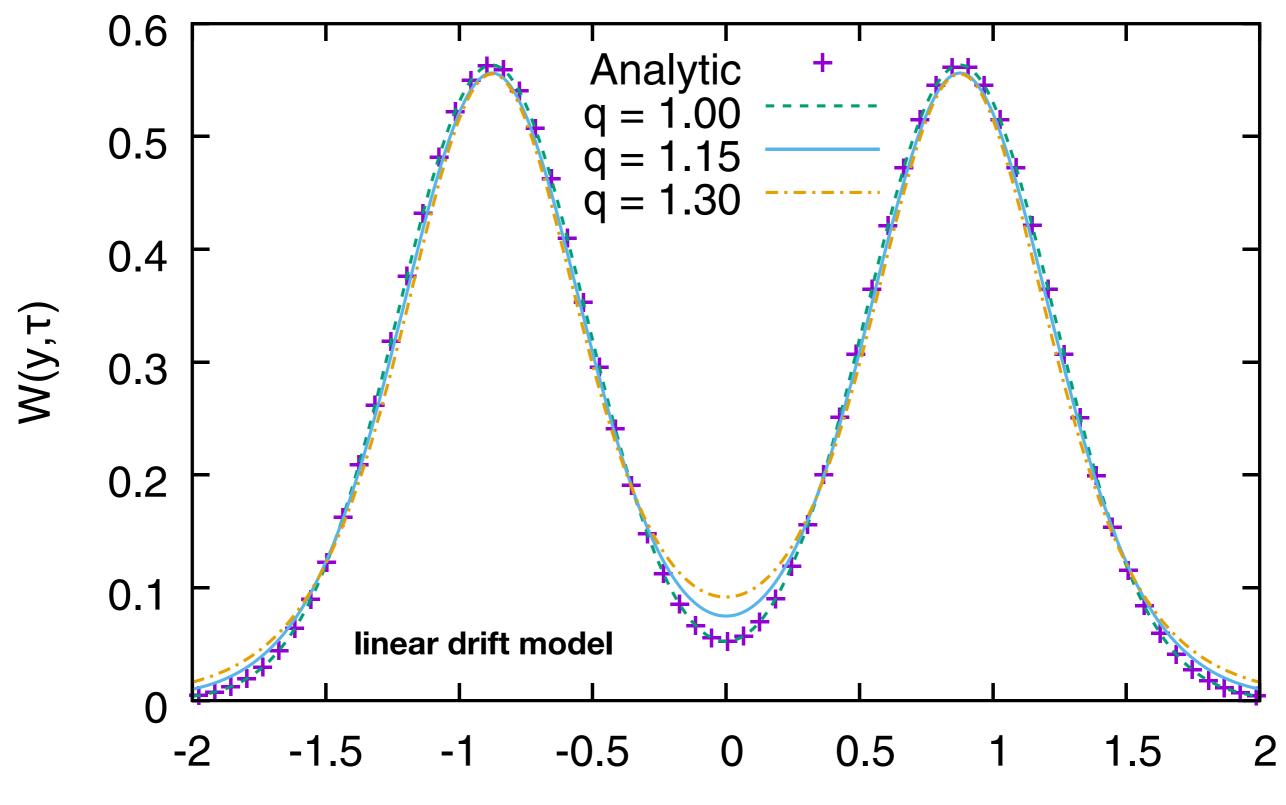
• Weak formulation (integral equation)

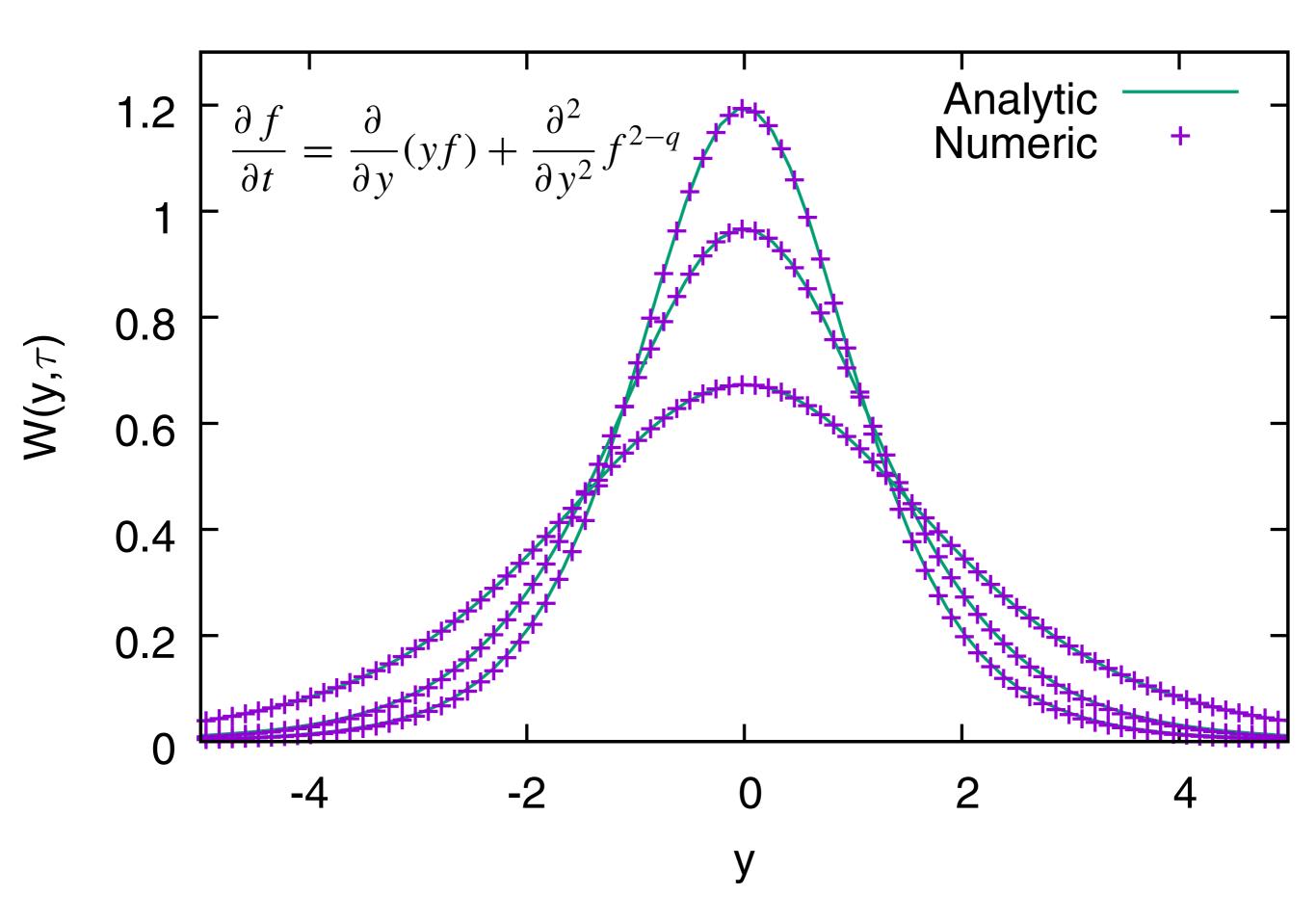
$$\begin{split} &\int_{\Omega} dy \left\{ g(y) \; \frac{\partial}{\partial y} \bigg[\sinh(y) \; f(y,t) + \gamma \; \frac{\partial}{\partial y} f(y,t)^{2-q} \bigg] \right\} \\ & \left[g(y) \bigg\{ \sinh(y) f(y,t) + \gamma \; \frac{\partial}{\partial y} f(y,t)^{2-q} \bigg\} \bigg] \bigg|_{\partial \Omega} \\ & - \int_{\Omega} dy \; \bigg\{ \frac{\partial g}{\partial y} \bigg[\sinh(y) f(y,t) + \gamma \; \frac{\partial}{\partial y} f(y,t)^{2-q} \bigg] \bigg\} \end{split}$$

• Time integration by Backward-Euler

$$\frac{\partial f(t_n)}{\partial t} = \frac{f(t_n) - f(t_{n-1})}{\Delta t} + O(\|\Delta t^2\|)$$

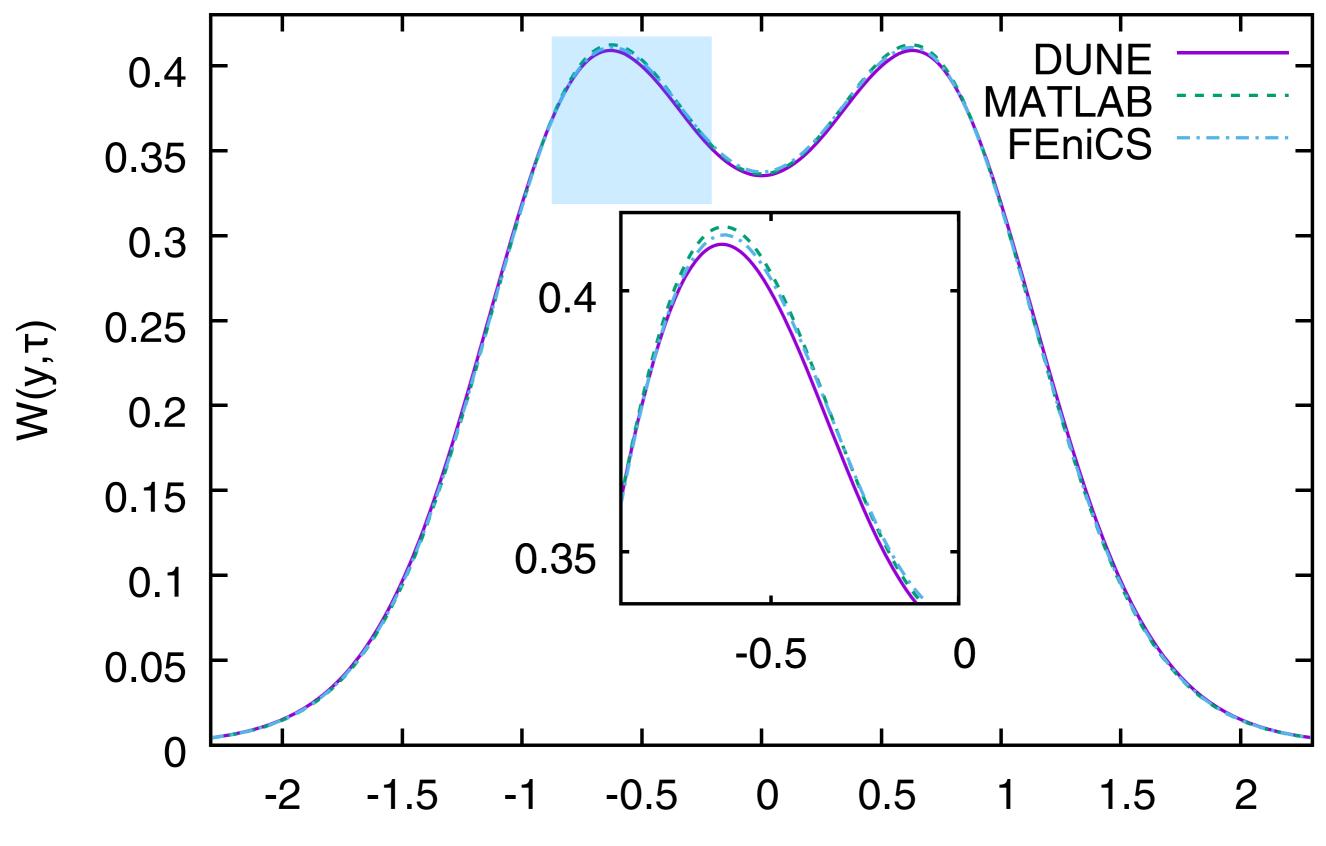
Numerical checks



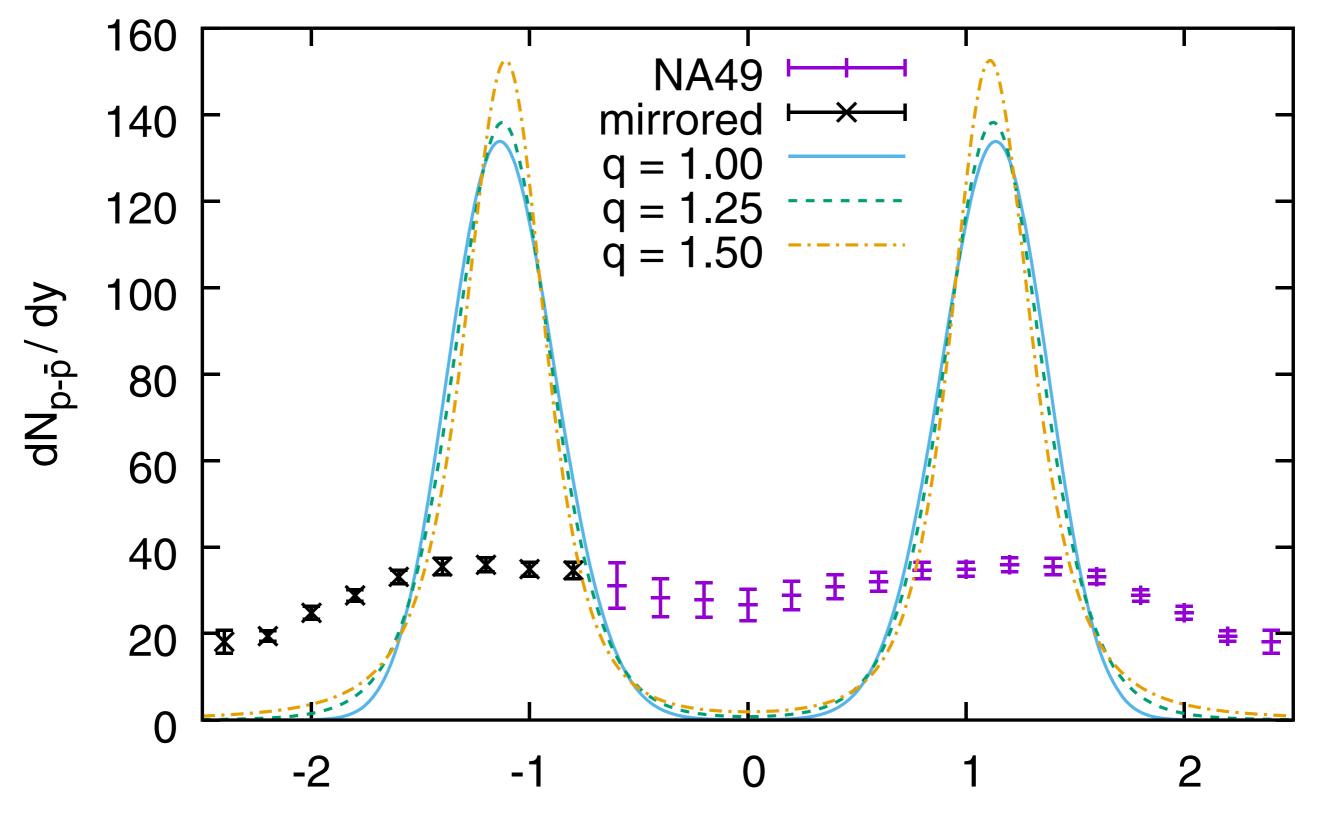


Borland, L., Pennini, F., Plastino, A. R., & Plastino, A. (1999). The nonlinear Fokker-Planck equation with state-dependent diffusion-a nonextensive maximum entropy approach. *The European Physical Journal B-Condensed Matter and Complex Systems*, *12*(2), 285-297.

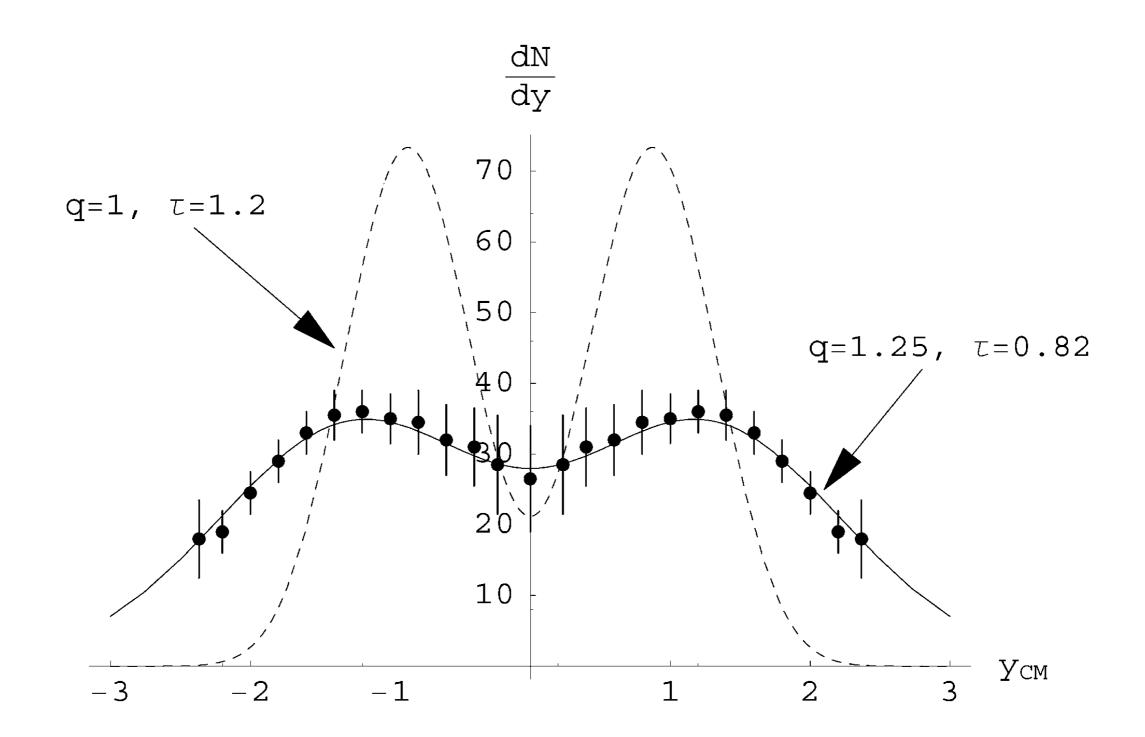
Method comparison



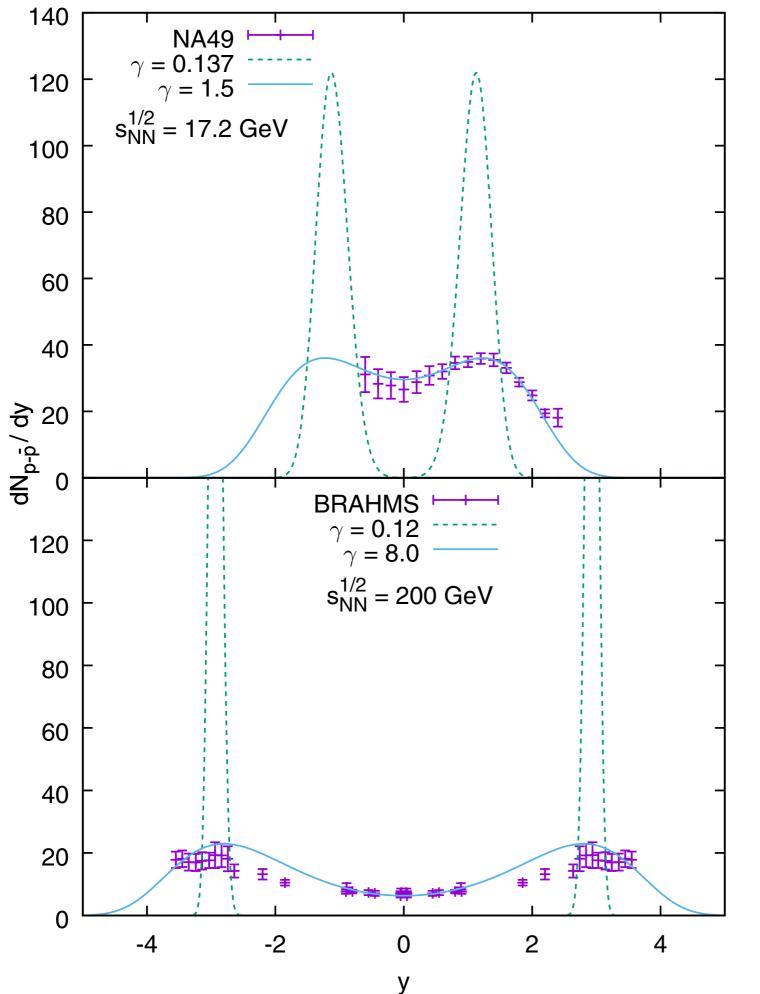
Result



≠ Lavagno's result



Lavagno, A. (2002). Anomalous diffusion in non-equilibrium relativistic heavy-ion rapidity spectra. *Physica A: Statistical Mechanics and its Applications*, 305(1), 238-241.



- Return to linear diffusion
- Adjust diffusion coefficient manually

Summary

- Measured rapidity spectra broader than anticipated
- Nonextensive statistics supposed to produce collective expansion
- Numerically solved the resulting nonlinear Fokker-Planck equation
- Effect is too weak to account for experimental data