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♀ suppression in pPb collisions

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Motivation

- pp collision QCD baseline
 - production mechanisms
- PbPb collision Hot Matter effects
 - ► Screening, Damping, Gluodissociation, Feed-Down reduction
- pPb collision Cold Nuclear Matter effects
 - Shadowing, Energy loss

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Basic Kinematics

- Lab System
- Lead Rest Frame

2 Cold Nuclear Matter Effects

- Collinear Factorization Theorem
- Nuclear Parton Distribution Function
- Coherent Parton Energy Loss

3 Y Results



Lab system

$$\bigcirc \xrightarrow{p_{\mathrm{P}} = \begin{pmatrix} E \\ p \end{pmatrix}} \xrightarrow{p_{\mathrm{N}} = \frac{Z}{A} \begin{pmatrix} E \\ -p \end{pmatrix}} < \xrightarrow{p_{\mathrm{N}} = \frac{Z}{A} \begin{pmatrix} E \\ -p \end{pmatrix}}$$

•
$$\sqrt{s_{\rm NN}} = \sqrt{(p_{\rm p} + p_{\rm N})^2} = \sqrt{s_{\rm pp}} \sqrt{\frac{Z}{A}} = 8.16 \ {\rm TeV}$$

•
$$\Delta y = \frac{1}{2} \ln \frac{E_{\text{cms}} + p_{\text{cms}}}{E_{\text{cms}} - p_{\text{cms}}} = \frac{1}{2} \ln \frac{A}{Z} = +0.465$$

•
$$y_{\text{beam}} = \pm \ln \frac{\sqrt{s_{\text{NN}}}}{m_{\text{p}}} = \pm 9.07$$

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Lead rest frame



•
$$s_{\rm NN} = 2E_{\rm p}m_{\rm p} \rightarrow E_{\rm p} = \frac{s_{\rm NN}}{2m_{\rm p}}$$

• $E_{\Upsilon} = M_{\perp} \cosh y' = M_{\perp} \cosh(y + y_{\text{beam}})$

• Suppression factor
$$R_{\text{pPb}}^{\Upsilon} = \frac{N_{\text{pPb}}^{\Upsilon}}{AN_{\text{pp}}^{\Upsilon}} = \frac{1}{A} \frac{\frac{\mathrm{d}\sigma_{\text{pPb}}}{\mathrm{d}y}}{\frac{\mathrm{d}\sigma_{\text{pp}}}{\mathrm{d}y}}$$

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4 Summary

Quarkonium Production mechanism



Collinear factorization theorem:

$$\sigma (\mathrm{pPb} \to \Upsilon + X) = \sum_{\mathrm{ij}} \int \mathrm{d}x_1 \mathrm{d}x_2 f_{\mathrm{i}}^{\mathrm{p}} (x_1) f_{\mathrm{j}}^{\mathrm{Pb}} (x_2) \sigma (\mathrm{ij} \to \Upsilon)$$

Parton Production

Color Evaporation Model

- Idea: Quark–Hadron duality
- $b\bar{b}$ pair will evaporate to a given color-singlet state via multiple soft gluon interaction
- Universal fraction to hadronize into Υ from a produced $b \hat{b}$ pair

$$\sigma\left(\mathrm{ij}\to\Upsilon\right) = F_{\Upsilon}\int_{4m^2}^{4M^2} \mathrm{d}s\;\hat{\sigma}\left(\mathrm{ij}\to b\bar{b}\right)$$

 $\bullet \,$ where $m(b)=4.18~{\rm GeV}$ and $M(B)=5.28~{\rm GeV}$

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- Above are leading order diagrams in QCD to calculate $b\hat{b}$ productions
- exclusive NLO heavy flavor calculations for $\frac{d\hat{\sigma}_{ij}}{dydp_{\perp}}$ are available as MNR code

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Partons



- $f_i^p(x, Q^2)$: probability distribution for finding a parton i depending on protons momentum fraction x and momentum transfer Q^2
- At higher energies there are not only valence quarks but also sea quarks and gluons

Parton Distribution Function



source: PHYSICAL REVIEW D 93, 033006 (2016)

Energy scale dependence given by DGLAP-equation

$$\frac{\partial f_i}{\partial \ln Q^2} = \frac{\alpha_s}{2\pi} \sum_j P_{i \leftarrow j} \otimes f_j$$

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Nuclear Modification



source: European Physical Journal C (2017) 77:163

- $f_i^{\text{Pb}}(x, Q^2) = R_i^{\text{Pb}}(x, Q^2) f_i^{\text{p}}(x, Q^2)$
- shadowing/depletion at small x, caused by gluon recombination
- anti-shadowing, an excess at $x \approx 0.1$ and EMC effect at higher x

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Coherent parton energy loss



source: Journal of High Energy Physics (2013) 2013: 122

- Assumption: The Υ forms outside of the nucleus
- Gluonbremsstrahlung of the partons and the produced heavy quarks

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coherent radiation



- Initial & Final state effect
- $q_{\perp} \gg l_{\perp}$ hard exchange; $l_{\perp}^2 = \bigtriangleup q_{\perp}^2$ nuclear momentum broadening
- Interference terms do not cancel in induced spectrum

$$\omega \frac{\mathrm{dI}}{\mathrm{d}\omega} = \frac{F_c \alpha_s}{\pi} \ln \left(1 + \frac{\hat{q}LE^2}{M_{\perp}^2 \omega^2} \right)$$

Energy shift

Modification of the cross section

$$\frac{1}{A}\frac{\mathrm{d}\sigma_{\mathrm{pPb}}}{\mathrm{d}E}(E,\sqrt{s}) = \int_{0}^{\varepsilon_{\mathrm{max}}} P(\varepsilon,E)\frac{\mathrm{d}\sigma_{\mathrm{pp}}}{\mathrm{d}E}(E+\varepsilon,\sqrt{s})\mathrm{d}\varepsilon$$

• $P(\epsilon)$: Quenching weight

$$P(\varepsilon) = \frac{\mathrm{dI}}{\mathrm{d}\omega}(\varepsilon) \exp\left(-\int_{\varepsilon}^{\infty} \mathrm{d}\omega \frac{\mathrm{dI}}{\mathrm{d}\omega}\right)$$

Average Energy Loss

$$\Delta E = \int \mathrm{d}\omega \; \omega \frac{\mathrm{d}\mathbf{I}}{\mathrm{d}\omega} \sim \alpha_s \frac{\sqrt{\hat{q}L}}{M_{\perp}} E$$

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source: Preliminary Results from the Quark Matter 2018 conference

- $\bullet\,$ Similar Υ suppression at forward and backward rapidity
- However significant forward suppression for J/ψ
- $\bullet~\Upsilon(2S)$ is slightly more suppressed than $\Upsilon(1S)$

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source: Preliminary Results from the Quark Matter 2018 conference

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Summary

Nuclear Modification of the Quarkonium Production

$$\sigma^{\rm pPb} = F_{\Upsilon} \sum_{\rm ij} \int_{4m^2}^{4M^2} ds \int dx_1 dx_2 f_{\rm i}^{\rm p}(x_1) f_{\rm j}^{\rm p}(x_2) \frac{R_j^{\rm Pb}(x_2)}{R_j^{\rm p}(x_2)} \hat{\sigma}_{\rm ij}$$

- Average Energy Loss $\triangle E \sim E$
- Both reproduce the stronger forward suppression, however the prediction slightly overestimate the results
- Color Evaporation Model fails to explain the stronger suppression of the excited state $\Upsilon(2S)$
- Does the Hot Matter Effect play a role?

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Thanks for your attention!

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