

Analytic model for the equilibration in finite Bose systems

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Derivation of the differential equation

- Hamiltonian consisting of free and interacting part

$$\hat{H} = \hat{H}_0 + \hat{H}_i$$

- von Neumann equation for weakly interacting bosons/fermions

$$\frac{\partial \langle \hat{n}_i \rangle}{\partial t} = - \left\langle \sum_{\gamma, \gamma'} |\gamma\rangle \langle \gamma| \Omega_{\gamma, \gamma'} \left(\langle \gamma | \hat{n}_i | \gamma \rangle - \langle \gamma' | \hat{n}_i | \gamma' \rangle \right) \right\rangle$$

- Transition amplitudes $\Omega_{\gamma, \gamma'}$

$$\Omega_{\gamma, \gamma'} = \left| \langle \gamma | \hat{H}_i | \gamma' \rangle \right|^2 G(\epsilon_\gamma, \epsilon_{\gamma'}) = \Omega_{\gamma', \gamma}$$

Weak $2 \rightarrow 2$ interaction

- Non-vanishing matrix elements for $|\gamma\rangle \rightarrow |\gamma'\rangle$

$$n'_i = \begin{cases} n_i & \text{for } i \neq i_1, i_2, i_3, i_4 \\ n_i - 1 & \text{for } i = i_1, i_2 \\ n_i + 1 & \text{for } i = i_3, i_4 \end{cases}$$

- Quantum Boltzmann equation for weakly interacting bosonic systems

$$\frac{\partial n_1}{\partial t} = \sum_{\epsilon_2, \epsilon_3, \epsilon_4}^{\infty} \left\langle V_{12,34}^2 \right\rangle G(\epsilon_1 + \epsilon_2, \epsilon_3 + \epsilon_4) \left[(1 + n_1)(1 + n_2)n_3n_4 - (1 + n_3)(1 + n_4)n_1n_2 \right]$$

Nonlinear diffusion equation I

- Master equation

$$\frac{\partial n_1}{\partial t} = (1 + n_1) \sum_{\epsilon_4} W_{4 \rightarrow 1} n_4 - n_1 \sum_{\epsilon_4} W_{1 \rightarrow 4} (1 + n_4)$$

- $\sum \rightarrow \int g_4 \, d\epsilon_4$ and Taylor expansion around ϵ_1

$$\frac{\partial n}{\partial t} = -\frac{\partial}{\partial \epsilon} \left[v n (1 + n) - n^2 \frac{\partial D}{\partial \epsilon} \right] + \frac{\partial^2}{\partial \epsilon^2} [Dn]$$

- Transport coefficients

$$D := \frac{1}{2} g_1 \int_0^\infty W(\epsilon_1, \epsilon_4) (\epsilon_1 - \epsilon_4)^2 \, d\epsilon_4$$

$$v := \frac{1}{g_1} \frac{\partial}{\partial \epsilon_1} (g_1 D)$$

Nonlinear diffusion equation II

- Assumption of constant transport coefficients

$$\frac{\partial n}{\partial t} = -v \frac{\partial}{\partial \epsilon} [n(1+n)] + D \frac{\partial^2 n}{\partial \epsilon^2}$$

- Bose-Einstein distribution with temperature $T = -\frac{D}{v}$ as static solution

$$n_{eq} = \frac{1}{\exp\left(\frac{\epsilon-\mu}{T}\right) - 1}$$

Analytic solution I

- Transformation to linear diffusion equation with partition function

$$n(\epsilon, t) = -\frac{D}{v} \frac{\partial}{\partial \epsilon} \ln(\mathcal{Z}) - \frac{1}{2} \quad \Rightarrow \quad \frac{\partial \mathcal{Z}}{\partial t} = D \frac{\partial^2 \mathcal{Z}}{\partial \epsilon^2}$$

- Partition function without boundary condition

$$\mathcal{Z}(\epsilon, t) = \int_{-\infty}^{+\infty} F(\epsilon, x, t) G(x) dx$$

- Partition function with boundary condition $\lim_{\epsilon \downarrow \mu} n(\epsilon, t) = \infty$

$$\mathcal{Z}(\epsilon, t) = \int_0^\infty \underbrace{(F(\epsilon - \mu, x, t) - F(\epsilon - \mu, -x, t))}_{\tilde{F}(\epsilon, x, t)} G(x + \mu) dx$$

Analytic solution II

- Green's function without boundary condition

$$F(\epsilon, x, t) = \exp\left(-\frac{(\epsilon - x)^2}{4Dt}\right)$$

- initial distribution (transformed)

$$G(x) = \exp\left(-\frac{1}{2D} \left(vx + 2v \int_0^x n_i(y) dy \right)\right)$$

- equivalent expression to avoid infinities

$$G(x) = \exp\left(-\frac{1}{2D} (vx + 2vA_i(x))\right)$$

Initial condition

- describe stochastic cooling
- initial condition:

$$n_i(\epsilon) = \frac{1}{\exp\left(\frac{\epsilon-\mu}{T}\right) - 1} \Theta(\epsilon_i - \epsilon)$$

- $T = -\frac{D}{v}$ → no description of cooling
- G - function

$$G(x) = \begin{cases} z^{-1} \exp\left(-\frac{vx}{2D}\right) - \exp\left(\frac{vx}{2D}\right) & \text{for } x < \epsilon_i \\ z^{-1} \exp\left(-\frac{vx}{2D}\right) - \exp\left(-\frac{vx}{2D} + \frac{v\epsilon_i}{D}\right) & \text{for } x > \epsilon_i \end{cases}$$

- fugacity: $z = \exp\left(\frac{\mu}{T}\right)$

Partition function

- Partition function for non-fixed chemical potential

$$\begin{aligned} \mathcal{Z} = & \sqrt{\pi D t} \exp\left(\frac{v^2 t}{4D}\right) \left[2z^{-1} \exp\left(-\frac{\epsilon v}{2D}\right) \right. \\ & - \exp\left(\frac{\epsilon v}{2D}\right) \operatorname{erfc}\left(\frac{\epsilon - \epsilon_i + tv}{\sqrt{4Dt}}\right) \\ & \left. - \exp\left(\frac{v\epsilon_i}{D} - \frac{\epsilon v}{2D}\right) \operatorname{erfc}\left(\frac{\epsilon_i - \epsilon + tv}{\sqrt{4Dt}}\right) \right] \end{aligned}$$

- Complementary error function:

$$\operatorname{erfc}(x) := 1 - \operatorname{erf}(x) := \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-t^2} dt$$

Particle distribution for non-fixed chemical potential

- Derivative of the error functions:

$$\exp\left(\frac{\epsilon v}{2D}\right) \frac{\partial}{\partial \epsilon} \operatorname{erfc}\left(\frac{\epsilon - \epsilon_i + tv}{\sqrt{4Dt}}\right) \\ + \exp\left(\frac{v\epsilon_i}{D} - \frac{\epsilon v}{2D}\right) \frac{\partial}{\partial \epsilon} \operatorname{erfc}\left(\frac{\epsilon_i - \epsilon + tv}{\sqrt{4Dt}}\right) = 0$$

- Particle distribution:

$$n(\epsilon, t) = \frac{1}{\exp\left(\frac{\epsilon - \mu}{T}\right) K(\epsilon, t) - 1}$$

- Auxiliary function $K(\epsilon, t)$:

$$K(\epsilon, t) = \frac{2 - \exp\left(\frac{\mu - \epsilon_i}{T}\right) \operatorname{erfc}\left(\frac{\epsilon_i - \epsilon + tv}{\sqrt{4Dt}}\right)}{\operatorname{erfc}\left(\frac{\epsilon - \epsilon_i + tv}{\sqrt{4Dt}}\right)}$$

Non-fixed μ - Solution for small times

parameters: $D = 8000 \text{ peV}^2\text{s}^{-1}$, $v = -1000 \text{ peVs}^{-1}$, $T = 8 \text{ peV}$,
 $\epsilon_i = 7 \text{ peV}$, $\mu = -0.68 \text{ peV}$

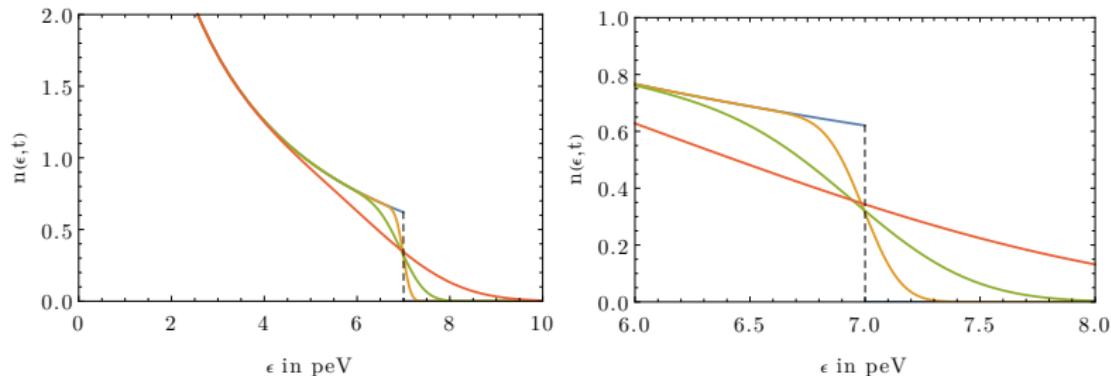


Figure: Solution for non-fixed μ at times 0 μs (blue), 1 μs (yellow), 10 μs (green) and 100 μs (red).

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Non-fixed μ - Large times

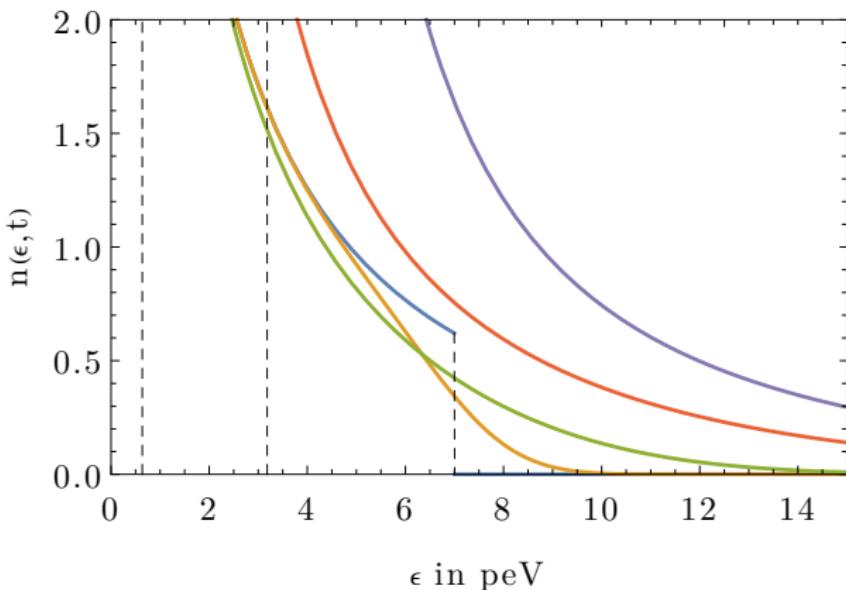


Figure: Solution for non-fixed μ at times $0\ \mu\text{s}$ (blue), $100\ \mu\text{s}$ (yellow), $1\ \text{ms}$ (green), $10\ \text{ms}$ (red) and at $t \rightarrow \infty$ (purple).

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Partition function

- Partition function for fixed chemical potential

$$\begin{aligned}\mathcal{Z}(\epsilon, t) = & \sqrt{4Dt} \exp\left(\frac{v^2 t}{4D} - \frac{\mu}{2T}\right) \left(\exp\left(\frac{\epsilon - \mu}{2T}\right) \Sigma_1(\epsilon, t) \right. \\ & \left. - \exp\left(\frac{\mu - \epsilon}{2T}\right) \Sigma_2(\epsilon, t) \right)\end{aligned}$$

- Auxiliary functions

$$\Sigma_1(\epsilon, t) = \operatorname{erfc}\left(\frac{2\mu - \epsilon_i - \epsilon + tv}{\sqrt{4Dt}}\right) - \exp\left(\frac{\mu - \epsilon_i}{T}\right) \operatorname{erfc}\left(\frac{\epsilon_i - \epsilon + tv}{\sqrt{4Dt}}\right)$$

$$\Sigma_2(\epsilon, t) = \operatorname{erfc}\left(\frac{\epsilon - \epsilon_i + tv}{\sqrt{4Dt}}\right) - \exp\left(\frac{\mu - \epsilon_i}{T}\right) \operatorname{erfc}\left(\frac{\epsilon - 2\mu + \epsilon_i + tv}{\sqrt{4Dt}}\right)$$

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Fixed μ - Solution

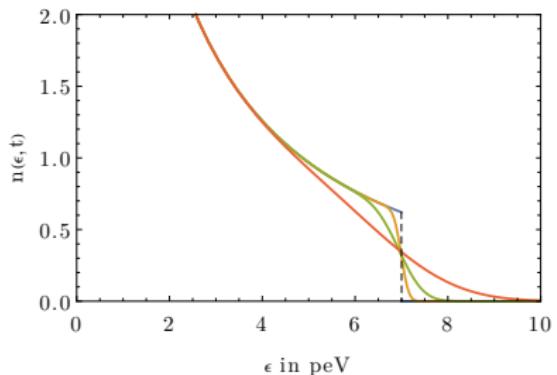


Figure: Solution for fixed μ at times 0 μ s (blue), 1 μ s (yellow), 10 μ s (green) and 100 μ s (red).

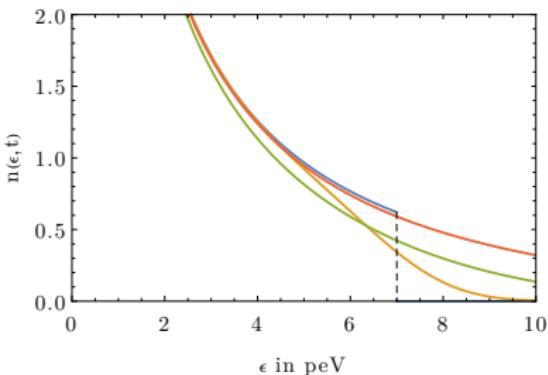


Figure: Solution for fixed μ at times 0 μ s (blue), 100 μ s (yellow), 1 ms (green) and 10 ms (red).

Partition function I - $G(x + \mu)$

- G - function for arbitrary initial temperature T_i

$$G(x + \mu) = \begin{cases} \exp\left(\frac{x-\mu}{2T_f}\right) \left(1 - \exp\left(-\frac{x}{T_i}\right)\right)^{\frac{T_i}{T_f}} & \text{for } x < \epsilon_i - \mu \\ \exp\left(\frac{x-\mu}{2T_f}\right) \left(1 - \exp\left(\frac{\mu-\epsilon_i}{T_i}\right)\right)^{\frac{T_i}{T_f}} & \text{for } x > \epsilon_i - \mu \end{cases}$$

- Separation into two parts:

$$\begin{aligned} \mathcal{Z}(\epsilon, t) &= \underbrace{\int_0^{\epsilon_i - \mu} \tilde{F}(\epsilon, x, t) G(x + \mu) dx}_{\mathcal{Z}_1(\epsilon, t)} \\ &\quad + \underbrace{\int_{\epsilon_i - \mu}^{\infty} \tilde{F}(\epsilon, x, t) G(x + \mu) dx}_{\mathcal{Z}_2(\epsilon, t)} \end{aligned}$$

Partition function II - $\mathcal{Z}_1(\epsilon, t)$

- Result for $\mathcal{Z}_1(\epsilon, t)$

$$\mathcal{Z}_1(\epsilon, t) = \sum_{k=0}^{\infty} \binom{\frac{T_i}{T_f}}{k} (-1)^k \exp\left(-\frac{\mu}{2T_f}\right) \underbrace{\int_0^{\epsilon_i - \mu} \tilde{F}(\epsilon, x, t) e^{\alpha x} dx}_{\mathcal{Z}_1^k(\epsilon, t)}$$

- Generalized binomial theorem and generalized binomial coefficients

$$(1+x)^s = \sum_{k=0}^{\infty} \binom{s}{k} x^k \quad , \quad \binom{s}{k} := \frac{s(s-1)\dots(s-k+1)}{k!}$$

- Auxiliary function $\alpha := \frac{1}{2T_f} - \frac{k}{T_i}$

Partition function III - $\mathcal{Z}_1^k(\epsilon, t)$

- Result for $\mathcal{Z}_1^k(\epsilon, t)$

$$\mathcal{Z}_1^k = \sqrt{\pi D t} e^{\alpha^2 D t} \left[e^{\alpha(\epsilon - \mu)} \Lambda_1^k(\epsilon, t) - e^{\alpha(\mu - \epsilon)} \Lambda_2^k(\epsilon, t) \right]$$

- Auxiliary functions

$$\Lambda_1^k(\epsilon, t) = \operatorname{erf} \left(\frac{\epsilon - \mu + 2D t \alpha}{\sqrt{4 D t}} \right) - \operatorname{erf} \left(\frac{\epsilon - \epsilon_i + 2D t \alpha}{\sqrt{4 D t}} \right)$$

$$\Lambda_2^k(\epsilon, t) = \operatorname{erf} \left(\frac{\mu - \epsilon + 2D t \alpha}{\sqrt{4 D t}} \right) - \operatorname{erf} \left(\frac{2\mu - \epsilon - \epsilon_i + 2D t \alpha}{\sqrt{4 D t}} \right)$$

Partition function IV - $\mathcal{Z}_2(\epsilon, t)$

- Result for $\mathcal{Z}_2(\epsilon, t)$

$$\begin{aligned}\mathcal{Z}_2 = & \left(1 - \exp\left(\frac{\mu - \epsilon_i}{T_i}\right)\right)^{\frac{T_i}{T_f}} \exp\left(-\frac{\mu}{2T_f}\right) \sqrt{\pi Dt} \exp\left(\frac{Dt}{4T_f^2}\right) \\ & \times \left[\exp\left(\frac{\epsilon - \mu}{2T_f}\right) \Lambda_3(\epsilon, t) - \exp\left(\frac{\mu - \epsilon}{2T_f}\right) \Lambda_4(\epsilon, t) \right]\end{aligned}$$

- Auxiliary functions

$$\Lambda_3(\epsilon, t) = \operatorname{erfc}\left(\frac{\epsilon_i - \epsilon + tv}{\sqrt{4Dt}}\right)$$

$$\Lambda_4(\epsilon, t) = \operatorname{erfc}\left(\frac{\epsilon - 2\mu + \epsilon_i + tv}{\sqrt{4Dt}}\right)$$

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Partition function V - Final result

$$\begin{aligned}\mathcal{Z} &= \sqrt{4Dt} \exp\left(-\frac{\mu}{2T_f}\right) \sum_{k=0}^{\infty} \binom{\frac{T_i}{T_f}}{k} (-1)^k \\ &\quad \times \left(e^{\alpha^2 Dt} \left[e^{\alpha(\epsilon-\mu)} \Lambda_1^k(\epsilon, t) - e^{\alpha(\mu-\epsilon)} \Lambda_2^k(\epsilon, t) \right] \right. \\ &\quad + \exp\left(\frac{(\mu-\epsilon_i)k}{T_i}\right) \exp\left(\frac{Dt}{4T_f^2}\right) \\ &\quad \left. \times \left[\exp\left(\frac{\epsilon-\mu}{2T_f}\right) \Lambda_3(\epsilon, t) - \exp\left(\frac{\mu-\epsilon}{2T_f}\right) \Lambda_4(\epsilon, t) \right] \right)\end{aligned}$$

Equilibration for $T_i = 20$ peV - Small times

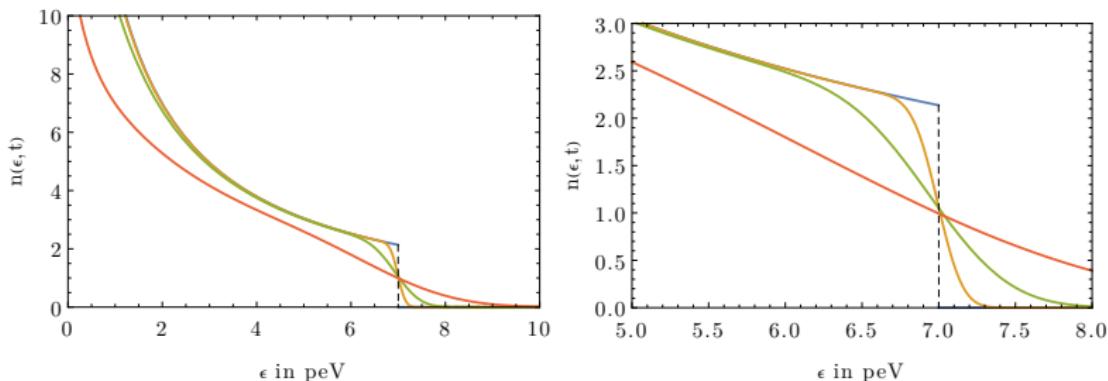


Figure: Equilibration starting with $T = 20$ peV at $0\text{ }\mu\text{s}$ (blue), $1\text{ }\mu\text{s}$ (yellow), $10\text{ }\mu\text{s}$ (green) and $100\text{ }\mu\text{s}$ (red).

Equilibration for $T_i = 20$ peV - Large Times

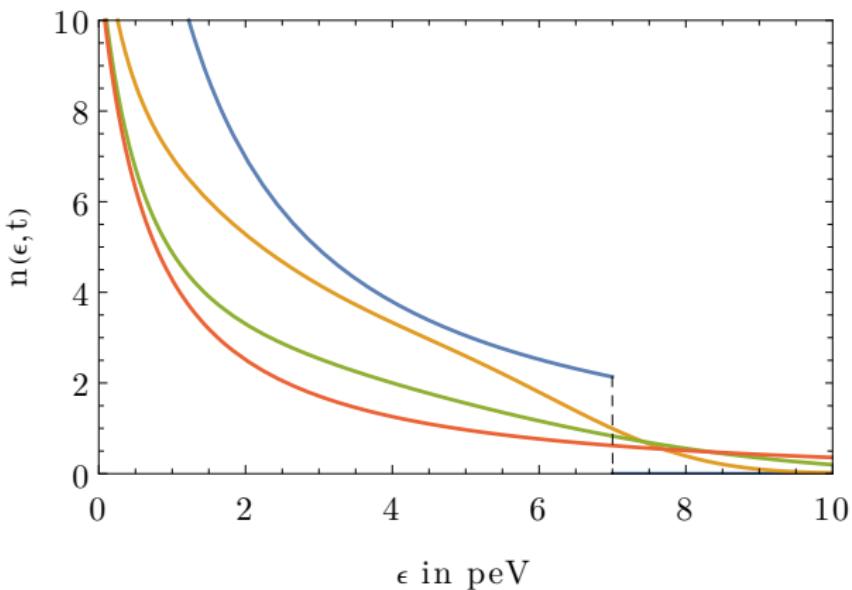


Figure: Equilibration starting with $T = 20$ peV at $0\ \mu\text{s}$ (blue), $100\ \mu\text{s}$ (yellow), $500\ \mu\text{s}$ (green) and $t \rightarrow \infty$ (red).

Number of particles

- Total particle number separated into thermal and condensate part

$$N = N_c(t) + N_t(t) = n(0, t) + \int_0^{\infty} g(\epsilon) n(\epsilon, t) d\epsilon$$

- Density of states for ideal uniform Bose gas

$$g(\epsilon) = g_0 \epsilon^k \quad , \quad g_0 = \frac{V}{4\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2}$$

- Particle number at $t = 0$

$$N = \frac{1}{z^{-1} - 1} + g_0 \int_0^{\epsilon_i} \frac{\sqrt{\epsilon}}{\exp\left(\frac{\epsilon - \mu}{T_i}\right) - 1} d\epsilon$$

Non-conserved particle number for constant μ

Parameters: $g_0 = 100 \text{ peV}^{-3/2}$, $k = \frac{1}{2}$

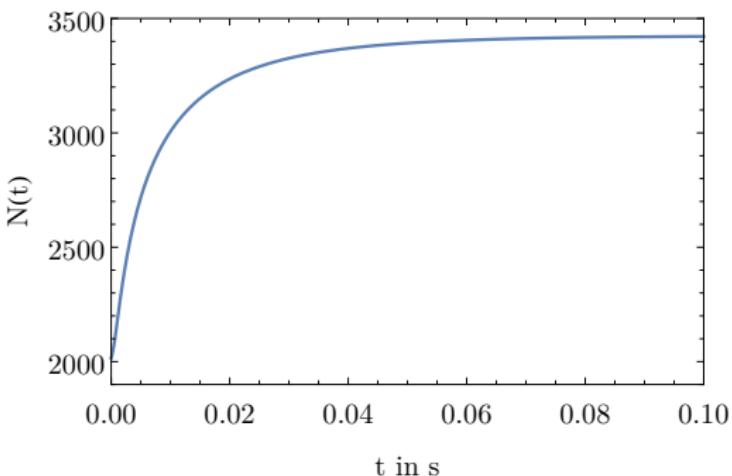


Figure: Particle number for fixed chemical potential $\mu = -0.68 \text{ peV}$.

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Time dependant chemical potential

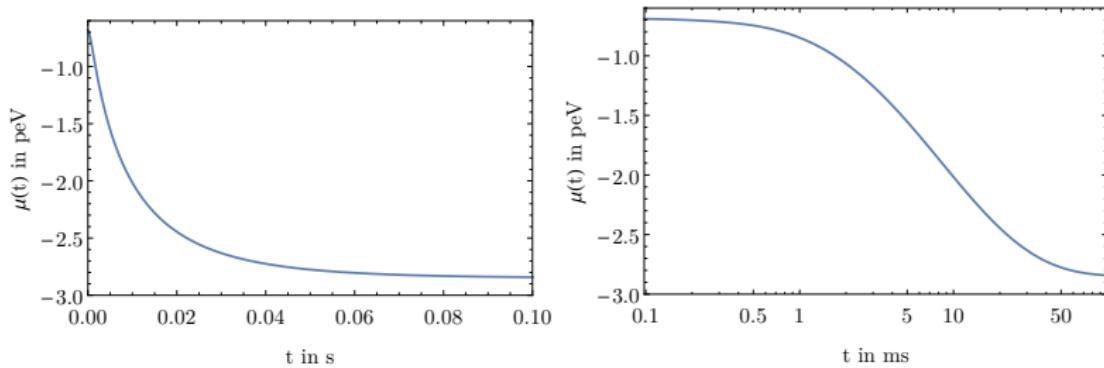


Figure: Time dependant chemical potential on linear time scale (left) and logarithmic time scale (right).

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Particle number conserving equilibration I

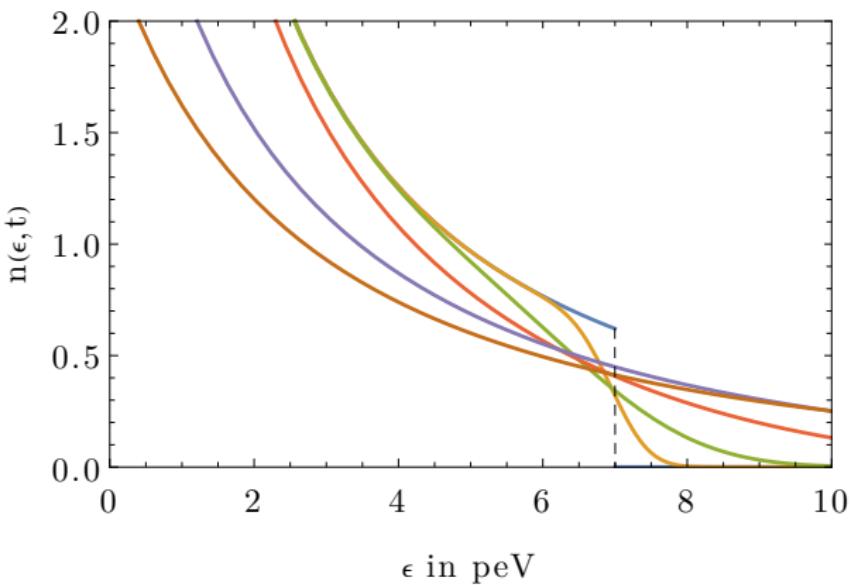


Figure: Equilibration for $T_i = T_f$ with time dependant, particle number conserving chemical potential at 0 μ s (blue), 10 μ s (yellow), 100 μ s (green), 1 ms (red), 10 ms (purple) and 100 ms.

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Arbitrary temperatures

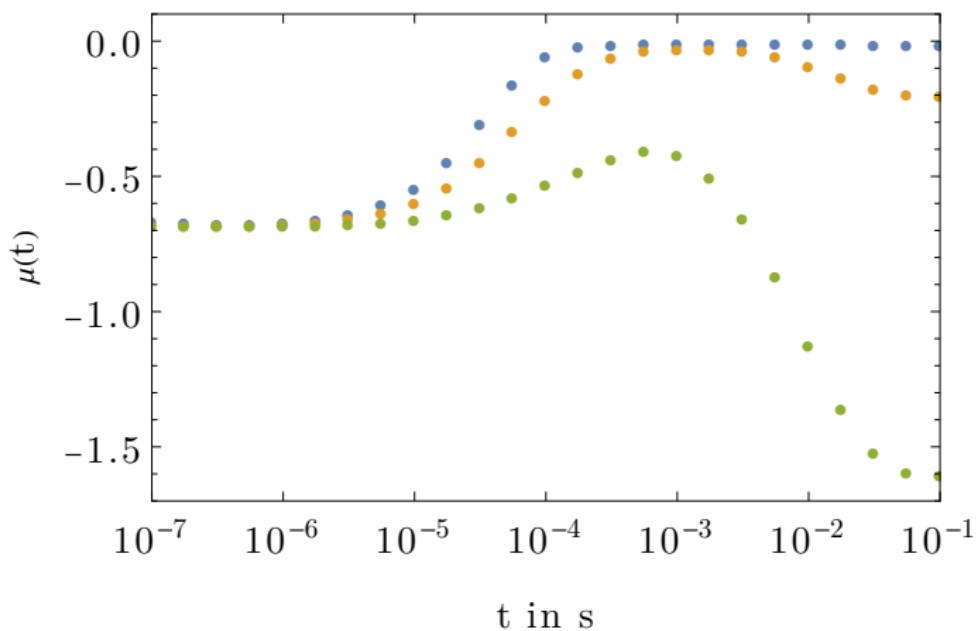


Figure: Time dependant chemical potential for initial temperature at 20 peV (blue), 15 peV (yellow) and 10 peV (green).

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Particle number conserving equilibration II

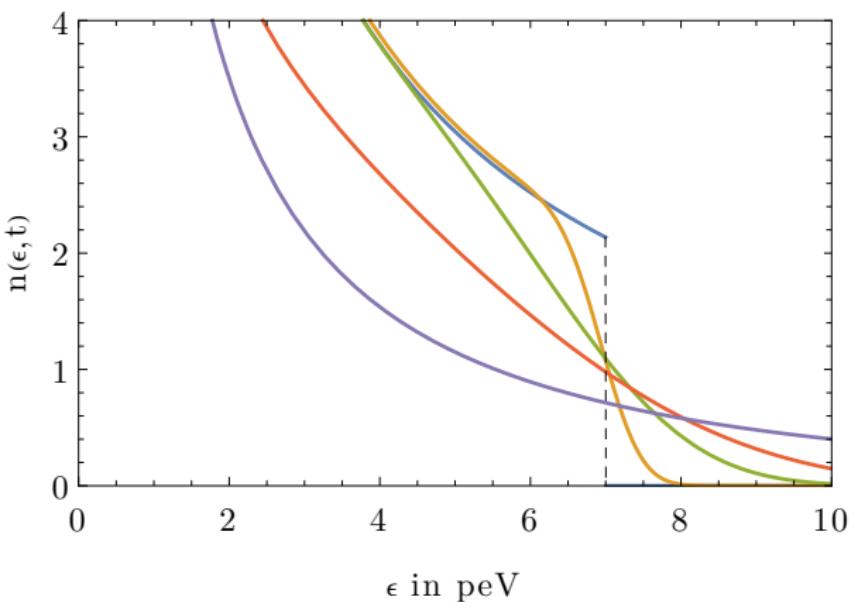


Figure: Equilibration for $T_i = 20$ peV with time dependant, particle number conserving chemical potential at $0\text{ }\mu\text{s}$ (blue), $10\text{ }\mu\text{s}$ (yellow), $100\text{ }\mu\text{s}$ (green), $316\text{ }\mu\text{s}$ (red) and $t \rightarrow \infty$ (purple).

Conclusion

- Converging solutions for the nonlinear diffusion equation
 - Analytic particle distributions
 - Two approaches, both not particle conserving
- Description of cooling, transition to different temperature
- Implementation of particle number conservation
 - Rise in condensed particles
 - No solutions anymore

References

-  G. Wolschin. "An exactly solvable model for equilibration in bosonic systems". In: *EPL (Europhysics Letters)* 123 (2018), p. 20009.
-  G. Wolschin. "Equilibration in finite Bose systems". In: *Physica A: Statistical Mechanics and its Applications* 499 (2018), p. 1.
-  P. Vasilopoulos and C. Van Vliet. "Quantum Boltzmann equations for binary interactions between identical particles". In: *Physica A: Statistical Mechanics and its Applications* 121 (1983), p. 617.
-  L. Pitaevskij and S. Stringari. *Bose-Einstein condensation*. International series of monographs on physics. Oxford: Clarendon Press, 2003.
-  *EqWorld - Linear Partial Differential Equations of Mathematical Physics*. <http://eqworld.ipmnet.ru/en/solutions/lpde.htm>. Accessed: 2019-06-27.
-  Otto Forster. *Analysis 1. Differential- und Integralrechnung einer Veränderlichen*. 10., überarbeitete und erweiterte Auflage. Wiesbaden: Vieweg+Teubner Verlag, 2011.